

INCLUSION OF DIPOLAR INTERACTIONS IN THE MATHEMATICAL MODELLING OF MAGNETIC HYPERTHERMIA THROUGH THE LANDAU-LIFSHITZ-GILBERT EQUATION PJ Cregg, Kieran Murphy & Mangolika Bhattacharya Waterford Institute of Technology, Ireland

Introduction

Magnetic Hyperthermia Treatment (MHT) continues to occupy clinicians, and experimental & theoretical magneticians alike. The many modelling approaches for non-interacting magnetic nanoparticles(MNPs) were outlined by Carrey *et al.* [1] at ICCSAMC 2014 in Dresden. Several authors have outlined the important role interparticle interactions are likely to play in MHT for closely spaced particles [1–6]. Modelling of MHT for non-interacting MNPs through the Landau-Lifshitz-Gilbert (LLG) equation was undertaken by Châtel *et al.* [7].

Our aim is to incorporate the dipole-dipole interactions in the LLG analysis. The effects of interactions on the frequency response of the heating mechanism are presented for two identical MNPs.

Effect of inter-particle distance



Relaxation mechanisms: Debye and Néel

• **Debye** relaxation: Brownian rotation of particle.

• Néel relaxation: Internal motion of magnetic moment - gyromagnetic. For a fixed particle - Debye blocked, Néel mechanism only - this can be described by LLG equation.

The Landau-Lifshitz-Gilbert (LLG) Equation

The LLG equation describes the average damped precessional motion of the magnetic moment, $\boldsymbol{\mu}$ (expressed as the normalised volume magnetisation, \mathbf{M} , where $\boldsymbol{\mu} = VM_s\mathbf{M}$) of an MNP in a magnetic field \mathbf{H} [7,8].

$$\dot{\mathbf{M}} = \left(\frac{\mu_0 \gamma}{1 + \alpha^2}\right) \mathbf{M} \times \mathbf{H} + \alpha \left(\frac{\mu_0 \gamma}{1 + \alpha^2}\right) \left(\mathbf{M} \times \mathbf{H}\right) \times \mathbf{M},$$

where

• α is the dimensionless damping parameter, defined by

 $\alpha = \eta \gamma M_s$

- η is the Gilbert dissipation constant,
- γ is the effective gyromagnetic ratio,
- μ_0 is permeability in vacuum,



(1)

(2)

d, distance between particle centres relative to particle radius.

Effect of particle centre alignment



 $\zeta,$ angle between external magnetic field and particle centres.

Conclusions

- Our results for the non-interacting case are consistent with ref. [7].
- Consistent with observations in ref. [6], the interparticle interactions are seen

- M_s is (volume) saturation magnetisation, and
- V is the volume of the MNP.

Inclusion of Interactions

The dipole-dipole interaction between MNPs can be included through the addition of the interaction field, (experienced by MNP 1 due to MNP 2), which for spherical MNPs of radius R, reduces to

 $\mathbf{H}_{\text{int}_1} = M_s \left(\frac{R}{|\mathbf{r}|}\right)^3 \left[\left(\mathbf{M}_2 \cdot \frac{\mathbf{r}}{|\mathbf{r}|}\right) \frac{\mathbf{r}}{|\mathbf{r}|} - \frac{1}{3}\mathbf{M}_2 \right]$

where \mathbf{r} is the vector between the MNP centres.

Calculation of Work done

We calculated the heating energy (per unit volume, per cycle), E, through the work done, i.e., damping force \times distance

$$E = \mu_0 \alpha M_s^2 \int_0^{2\pi/\omega} \dot{\mathbf{M}} \cdot d\mathbf{M} = \mu_0 \alpha M_s^2 \int_0^{2\pi/\omega} |\dot{\mathbf{M}}|^2 dt$$
(3)

This can be shown to be analytically equivalent to that of Châtel *et al.* [7] where ω is the angular frequency of **H**,

to hinder the heating mechanism. As expected the interaction effects fall off with $\sim |\mathbf{r}|^3$. Thus, interaction effects predominate when interparticle distance is $\leq \sqrt[3]{\frac{M_s}{H}}$.

Future Work

More complex arrangements can be considered. These include likely physical arrangements such as: • Chains • Cluster • Equi-spaced • Random

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Behaviour of two MNPs

The heating per cycle, E and the consequent Specific Absorption Rate (SAR) (taken to be the per unit volume value found from SAR = Ef, for frequency f) can be calculated over a range of frequencies.

The effect of dipole-dipole interactions is investigated for different MNP orientations.

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